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THEORETICAL STUDYING SPECTRA OF YTTERBIUM ATOM ON THE BASIS OF RELATIVISTIC MANY-BODY PERTURBATION THEORY: RYDBERG RESONANCES

Theoretical studying the Rydberg autoionization resonances in spectra of the lanthanides atoms (ytterbium) is carried out within the relativistic many-body perturbation theory and generalized relativistic energy approach (Gell-Mann and Low S-matrix formalism). The zeroth approximation of the relativistic perturbation theory is provided by the optimized Dirac-Fock and Dirac-Kohn-Sham ones. Optimization has been fulfilled by means of introduction of the parameter to the Fock and Kohn-Sham exchange potentials and further minimization of the gauge-non-invariant contributions into radiation width of atomic levels with using relativistic orbital bases, generated by the corresponding zeroth approximation Hamiltonians. The accurate theoretical results on the autoionization $4f^{13} [^2F_{7/2}] 6s^2np[5/2]_2$, $4f^{13} [^2F_{7/2}] 6s^2nf[5/2]_2$ resonances energies and widths are presented and compared with experimental data, obtained on the basis of the laser polarization spectroscopy method.

1. Introduction

This paper goes on our work on theoretical studying spectra and spectroscopic parameters for heavy atoms, namely, lanthanides atoms (see, for example [1-56]). Let us remind that an investigation of spectra, optical and spectral, radiative and autoionization characteristics for heavy elements atoms and multicharged ions is traditionally of a great interest for further development quantum atomic optics and atomic spectroscopy and different applications in plasma chemistry, astrophysics, laser physics etc. (see Refs. [1-10]).

Different atomic spectroscopy methods have been used in studying radiative and autoionization characteristics of atomic systems. The well known classical multi-configuration Hartree-Fock method allowed to get a great number of the useful spectral information about light and not heavy atomic systems. The multi-configuration Dirac-Fock method is the most reliable version of calculation for multielectron systems with a large nuclear charge. In these calculations the one- and two-particle relativistic and important exchange-correlation corrections are taken into account (see Refs. [1] and Refs. therein). However, one should remember about very complicated

structure of spectra of the lanthanides atoms and necessity of correct accounting the different correlation effects such as polarization interaction of the valent quasiparticles and their mutual screening, iterations of a mass operator etc.). The known method of the model relativistic many-body perturbation theory has been earlier effectively applied to computing spectra of low-lying states for some lanthanides atoms too [2,3] (see also [2-6]). We use an analogous version of the perturbation theory (PT) to study the Rydberg states characteristics, however, the zeroth approximation is generated by the Dirac-Fock model. In Refs. [7-10] the similar version of the perturbation theory has been used with using the Dirac-Kohn-Sham zeroth approximation. This method is actively used in solving many tasks of quantum, atomic and nuclear physics [57-87]. Here we present the results of computing the Rydberg Yb $4f^{13} [^2F_{7/2}] 6s^2np[5/2]_2$, $4f^{13} [^2F_{7/2}] 6s^2nf[5/2]_2$ states energies and widths within both approaches and compare theoretical data with some experimental laser polarization spectroscopy method data [22,23]. All calculations are performed with using Superatom package (see for example, 2-24]).

2. Advanced relativistic many-body perturbation theory and energy approach

As the method of computing is earlier presented in details, here we are limited only by the key topics. A model relativistic energy approach in a case of the multielectron atom has been proposed by Ivanov-Ivanova et al [2-4] and its generalized gauge-invariant version is developed in Refs. [5,6,11,2]. The approach is based on the Gell-Mann and Low S-matrix formalism and the relativistic many-body PT with using the optimized one-quasiparticle representation and an accurate account of the relativistic and exchange-correlation effects. In the relativistic case the Gell-Mann and Low formula expressed an energy shift ΔE through the QED scattering matrix including the interaction with as the photon vacuum field as the laser field. The wave function zeroth basis is found from the Dirac equation with a potential, which includes ab initio optimized model (Ivanov-Ivanova type [6]) potential or DF potentials, the electric potential of a nucleus (the Gaussian form of the charge distribution in a nucleus is usually used by us) [4]. The correlation corrections of the PT second and higher orders are taken into account by means of using the polarization and screening potentials (from Refs. [10-16]).

Generally speaking, the majority of complex atomic systems possess a dense energy spectrum of interacting states with essentially relativistic properties. In the theory of the non-relativistic atom a convenient field procedure is known for calculating the energy shifts ΔE of degenerate states. This procedure is connected with the secular matrix M diagonalization [2,11,12]. In constructing M , the Gell-Mann and Low adiabatic formula for ΔE is used. In contrast to the non-relativistic case, the secular matrix elements are already complex in the second order of the electro-dynamical PT (first order of the interelectron interaction). Their imaginary part of ΔE is connected with the radiation decay (radiation) possibility. In this approach, the whole calculation of the energies and decay probabilities of a non-degenerate excited state is reduced to the calculation and diagonalization of the complex matrix M . In the papers of different authors, the $\text{Re}\Delta E$ calculation procedure has been generalized for the case of nearly degenerate states, whose levels form a

more or less compact group. One of these variants has been previously [7-12] introduced: for a system with a dense energy spectrum, a group of nearly degenerate states is extracted and their matrix M is calculated and diagonalized. If the states are well separated in energy, the matrix M reduces to one term, equal to ΔE . The non-relativistic secular matrix elements are expanded in a PT series for the interelectron interaction. The complex secular matrix M is represented in the form [3,4,11]:

$$M = M^{(0)} + M^{(1)} + M^{(2)} + M^{(3)}, \quad (1)$$

where $M^{(0)}$ is the contribution of the vacuum diagrams of all order of PT, and $M^{(1)}$, $M^{(2)}$, $M^{(3)}$ those of the one-, two- and three- quasiparticle diagrams respectively. $M^{(0)}$ is a real matrix, proportional to the unit matrix. It determines only the general level shift. We have assumed $M^{(0)} = 0$. The diagonal matrix $M^{(1)}$ can be presented as a sum of the independent one-quasiparticle contributions. For simple systems (such as alkali atoms and ions) the one-quasiparticle energies can be taken from the experiment. Substituting these quantities into (1) one could have summarized all the contributions of the one -quasiparticle diagrams of all orders of the formally exact QED PT. However, the necessary experimental quantities are not often available. The first two order corrections to $\text{Re}M^{(2)}$ have been analyzed previously [4] using Feynman diagrams [11]. The contributions of the first-order diagrams have been completely calculated. In the second order, there are two kinds of diagrams: polarization and ladder ones. The polarization diagrams take into account the quasiparticle interaction through the polarizable core, and the ladder diagrams account for the immediate quasiparticle interaction [11-20]. Some of the ladder diagram contributions as well as some of the three-quasiparticle diagram contributions in all PT orders have the same angular symmetry as the two-quasiparticle diagram contributions of the first order. These contributions have been summarized by a modification of the central potential, which must now include the screening (anti-screening) of the core potential of each particle by the two others. The additional potential modifies the one-quasiparticle orbitals and energies. Then the secular matrix is as follows:

$$M \rightarrow \tilde{M}^{(1)} + \tilde{M}^{(2)}, \quad (2)$$

where $\tilde{M}^{(1)}$ is the modified one-quasiparticle matrix (diagonal), and $\tilde{M}^{(2)}$ the modified two-quasiparticle one. $\tilde{M}^{(1)}$ is calculated by substituting the modified one-quasiparticle energies, and $\tilde{M}^{(2)}$ by means of the first PT order formulae for $M^{(2)}$, putting the modified radial functions of the one-quasiparticle states in the radial integrals.

Let us remind that in the QED theory, the photon propagator $D(12)$ plays the role of this interaction. Naturally the analytical form of $D(12)$ depends on the gauge, in which the electrodynamic potentials are written. Interelectron interaction operator with accounting for the Breit interaction has been taken as follows:

$$V(r_i r_j) = \exp(i\omega r_{ij}) \cdot \frac{(1 - \alpha_i \alpha_j)}{r_{ij}}, \quad (3)$$

where, as usually, α_i are the Dirac matrices. In general, the results of all approximate calculations depended on the gauge. Naturally the correct result must be gauge-invariant. The gauge dependence of the amplitudes of the photo processes in the approximate calculations is a well known fact and is in details investigated by Grant, Armstrong, Aymar and Luc-Koenig, Glushkov-Ivanov et al (see [11] and numerous Refs. therein). Grant has investigated the gauge connection with the limiting non-relativistic form of the transition operator and has formulated the conditions for approximate functions of the states, in which the amplitudes of the photo processes are gauge invariant [1]. These results remain true in the energy approach because the final formulae for the probabilities coincide in both approaches. Glushkov-Ivanov have developed a new relativistic gauge-conserved version of the energy approach [6]. Here we applied this approach for generating the optimized relativistic orbitals basis in the zeroth approximation of the many-body PT [7-10]. Optimization has been fulfilled by means of introduction of the parameter to the Fock and Kohn-Sham exchange potentials and further minimization of the gauge-non-invariant contributions into radiation width of atomic levels with using relativistic orbital bases, generated by the corresponding zeroth approximation Hamiltonians. Below we will be interested by studying the spectra of

the autoionization resonances in the ytterbium atom and calculating their energies and widths. The excited states of the ytterbium atom can be treated as the states with two-quasiparticles above the electron core $[\text{Xe}]4f^{14}$. Within the standard energy approach [8-11] the autoionization width is determined by the square of an electron interaction (3) matrix element. The real part of the electron interaction matrix element is determined using expansion in terms of Bessel functions [17-19,26]; the Coulomb part Q_λ^{Coul} is expressed in terms of the standard radial integrals and angular coefficients. The Breit part of Q is defined in the similar way as above, but the contribution of our interest is a real part. The Breit interaction is known to change considerably the autoionization decay dynamics in some cases (see, for example, Refs. [3,11]). Determination of the radiation decay probabilities (oscillator strengths) results to calculating the imaginary matrix elements of the interaction (3). The calculation of radial integrals $\text{Re}R_\lambda(1243)$ is reduced to the solution of a system of differential equations according to the Ivanova-Ivanov method [26]. The system of differential equations includes also equations for functions $f/r^{|\lambda|-1}$, $g/r^{|\lambda|-1}$, $Z_\lambda^{(1)}$, $Z_\lambda^{(2)}$. The formulas for the autoionization decay probability include the radial integrals $R_\alpha(\alpha k \gamma \beta)$, where one of the functions describes electron in the continuum state. The correctly normalized function should have the following regular asymptotic at $r \rightarrow 0$ (look details in Refs. [13-19]). Other details can be found in Refs. [6-11,13-19].

3. Some illustration results and conclusion

In table 1 we present the experimental data (Jong-Hoon et al [87,88]) and theoretical results (Th1-PT with the Dirac-Kohn-Sham zeroth approximation [7,8]; Th2 – this work) for energies and widths of the excited (autoionization) states of the $4f^{13} [{}^2F_{7/2}]6s^2np[5/2]_2$ and $4f^{13} [{}^2F_{7/2}] 6s^2nf [5/2]_2$ states (because of excitation of the $4f$ shell).

In table 2 we present the predictions of this work regarding the energies and widths of the excited (autoionization) states $4f^{13} [{}^2F_{7/2}] 6s^2nf [5/2]_2$ states. As it has been noted in [5], the attention is drawn to the smallness of the resonance widths, the cause of which in the literature is not clear. In

our opinion, it is related to the complex energetic structure of the 4f-shell atoms, as a result of causing several unusual physics of autoionization resonances and their decay mechanisms, especially in comparison with the conventional standards spectroscopy (for He, inert gases, alkali atoms)

Table 1
Energies E (cm^{-1}), widths Γ (cm^{-1}) of the $4f^{13} [^2F_{7/2}]6s^2np[5/2]_2$ states in YbI: Th1- PT with Dirac-Kohn-Sham zeroth approximation; Th2- Th1- PT with Dirac-Fock zeroth approximation (this work)

n	Exp. E_{exp}	Exp. Γ_{exp}	Th1. E	Th1. Γ	Th2. E	Th2. Γ
12	70120.5	1.5	70121	1.7	70123	1.6
15	70914.8	1.2	70916	1.4	70917	1.3
20	71428.1	0.6	71429	0.7	71430	0.6
25	71612.5	1.3	71611	1.5	71612	1.4
26	71633.3	0.6	71631	0.8	71633	0.7
30	71698.8	0.5	71697	0.7	71699	0.6
46	-	-	-	-	71798	0.3

It is important to note that the both perturbation theory versions with the Dirac-Fock and Dirac-Kohn-Sham zeroth approximations provide a physically reasonable agreement with experiment, however, more exact data are provided by the optimized Dirac-Fock-like theory.

Table 2
Energies E (cm^{-1}), widths Γ (cm^{-1}) of the $4f^{13} [^2F_{7/2}] 6s^2 nf [5/2]_2$ states (predictions of this work)

n	Th. E	Th. Γ
12	70966	0.6
13	71109	0.4
15	71314	1.5
20	71562	0.8
25	71674	0.7
26	71690	0.6
30	71735	0.4
46	71814	0.2

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THEORETICAL STUDYING SPECTRA OF YTTERBIUM ATOM ON THE BASIS OF RELATIVISTIC MANY-BODY PERTURBATION THEORY: RYDBERG RESONANCES

Summary

Theoretical studying the Rydberg autoionization resonances in spectra of the lanthanides atoms (ytterbium) is carried out within the relativistic many-body perturbation theory and generalized relativistic energy approach (Gell-Mann and Low S-matrix formalism). The zeroth approximation of the relativistic perturbation theory is provided by the optimized Dirac-Fock and Dirac-Kohn-Sham ones. Optimization has been fulfilled by means of introduction of the parameter to the Fock and Kohn-Sham exchange potentials and further minimization of the gauge-non-invariant contributions into radiation width of atomic levels with using relativistic orbital bases, generated by the corresponding zeroth approximation Hamiltonians. The accurate theoretical results on the autoionization $4f^{13} [^2F_{7/2}] 6s^2 np [5/2]_2$, $4f^{13} [^2F_{7/2}] 6s^2 nf [5/2]_2$ resonances energies and widths are presented and compared with experimental data, obtained on the basis of the laser polarization spectroscopy method.

Keywords: Relativistic perturbation theory, resonances energies and widths, optimized zeroth approximation

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**ТЕОРЕТИЧЕСКОЕ ИЗУЧЕНИЕ СПЕКТРА АТОМА ИТТЕРБИЯ НА ОСНОВЕ
РЕЛЯТИВИСТСКОЙ МНОГОЧАСТИЧНОЙ ТЕОРИИ ВОЗМУЩЕНИЙ:
РИДБЕРГОВЫ РЕЗОНАНСЫ**

Резюме

В рамках релятивистской многочастичной теории возмущений и обобщенного релятивистского энергетического подхода проведено теоретическое изучение характеристик ридберговских автоионизационных резонансов в спектрах атомов лантанидов (иттербия). В качестве нулевого приближения релятивистской теории возмущений выбраны оптимизированные приближения Дирака-Фока и Дирака-Кона-Шэма. Оптимизация выполнена путем введения параметра в обменные потенциалы Фока и Кона-Шэма и дальнейшей минимизацией калибровочно-неинвариантных вкладов в радиационные ширины атомных уровней с использованием релятивистских орбитальных базисов, сгенерированных соответствующими гамильтонианами нулевого приближения. Представлены аккуратные теоретические данные по энергиям и ширинам автоионизационных $4f^{13} [{}^2F_{7/2}] 6s^2 np [5/2]_2$, $4f^{13} [{}^2F_{7/2}] 6s^2 nf [5/2]_2$ резонансов и проведено сравнение с экспериментальными данными, полученными на основе метода лазерной поляризационной спектроскопии.

Ключевые слова: Релятивистская теория возмущений, энергии и ширины резонансов, оптимизированное нулевое приближение

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**ТЕОРЕТИЧНЕ ВИВЧЕННЯ СПЕКТРА АТОМА ІТЕРБІЮ НА ОСНОВІ
РЕЛЯТИВІСТСЬКОЇ БАГАТОЧАСТИНКОВОЇ ТЕОРІЇ ЗБУРЕНЬ: РІДБЕРГОВІ
РЕЗОНАНСИ**

Резюме

В рамках релятивістської багаточастинкової теорії збурень і узагальненого релятивістського енергетичного підходу проведено теоретичне вивчення характеристик рідбергівських автоіонізаційних резонансів в спектрах атомів лантанідів (ітербію). В якості нульового наближення релятивістської теорії збурень обрані оптимізовані наближення Дірака-Фока і Дірака-Кона-Шема. Оптимізація виконана шляхом введення параметра в обмінні потенціали Фока і Кона-Шема і подальшої мінімізації калібрувальних неінваріантних вкладів в радіаційні ширини атомних рівнів з використанням релятивістських орбітальних базисів, згенерованими відповідними гамільтоніанами нульового наближення. Представлені акуратні теоретичні дані по енергіях і ширинам автоіонізаційних $4f^{13} [{}^2F_{7/2}] 6s^2 np [5/2]_2$, $4f^{13} [{}^2F_{7/2}] 6s^2 nf [5/2]_2$ резонансів і проведено порівняння з експериментальними даними, отриманими на основі методу лазерної поляризаційної спектроскопії.

Ключові слова: Релятивістська теорія збурень, енергії і ширини резонансів, оптимізоване нульове наближення