

THE INVARIANT SPLITTING FORMALISM IN GENERAL RELATIVITY

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ABSTRACT. The invariant splitting method of a pseudo-Riemannian manifold on arbitrary number of sections has been constructed in the present work.

Key words: general relativity, splitting, tangent bundle, the Gauss-Codacci-Ricci's equations.

Our technique is based on introduction of projection operators $\{H^\alpha\}$, inducing splitting of the bundle $T(M)$ and geometrical objects defined on it. The linear operators $\{H^\alpha\}$ (the tensors of (1,1) type) are given by:

$$\begin{aligned} H^\alpha \cdot H^\beta &= \delta^{\alpha\beta} \\ \sum_\alpha H^\alpha &= I \end{aligned} \quad (1)$$

where $\alpha, \beta = 1, 2, \dots, k$ are indices of sections and I is the unit operator. Any vector \vec{X} and one-form ω are projected into corresponding sections of the tangent and cotangent bundles by these projection operators $\{H^\alpha\}$

$$\begin{aligned} H^\alpha \cdot \vec{X}^\alpha &\equiv H^\alpha(\vec{X}) = \vec{X}^\alpha, \vec{X}^\alpha \in \Sigma^\alpha \\ H^\alpha \cdot \omega &\equiv H^\alpha(\omega) = \omega^\alpha, \omega^\alpha \in \Sigma^{*\alpha} \\ T(M) &= \Sigma^1 \oplus \Sigma^2 \oplus \dots \oplus \Sigma^k \\ \dim T(M) &= \sum_{i=1}^k \dim \Sigma^i, \\ \vec{X} &= \sum_{\alpha=1}^k \vec{X}^\alpha, \\ \Sigma^\alpha &\equiv \text{Im} H^\alpha \end{aligned} \quad (2)$$

so that for any metric g we have

$$g = \sum_{\alpha=1}^k g^\alpha \quad (3)$$

where g^α is the metric induced on the section Σ^α . Due to conditions $H^\alpha \cdot H^\beta = 0$ and $g(H^\alpha \cdot \vec{X}, \vec{Y}) = g(\vec{X} \cdot H^\alpha, \vec{Y})$ it follows that

$$\vec{X}^\alpha \cdot \vec{Y}^\beta = 0, g^\alpha(\vec{X}^\beta, \cdot) = 0 \quad (4)$$

The Levi-Civita connection ∇ is splitted on k^3 components

$$\nabla_{\vec{X}} \vec{Y} = \sum_{\alpha, \beta, \gamma} \nabla_{\vec{X}^\alpha}^\gamma \vec{Y}^\beta, \quad \alpha, \beta, \gamma = 1, 2, \dots, k \quad (5)$$

and the components

$$\nabla_{\vec{X}^\beta}^\alpha \vec{Y}^\beta = H^\alpha(\nabla_{\vec{X}^\beta} \vec{Y}^\beta) \equiv -B(\vec{X}^\beta, \vec{Y}^\beta)$$

determine the exterior unholonomity tensor of the section Σ^β , while the components

$$\nabla_{\vec{X}^\beta}^\alpha \vec{Y}^\beta = H^\alpha(\nabla_{\vec{X}^\beta} \vec{Y}^\gamma) \equiv -Q^\alpha(\vec{X}^\beta, \vec{Y}^\gamma)$$

determine the generalized Ricci's coefficients of rotation.

Substituting all the connections in the invariant definition of the curvature tensor

$$R(\vec{X}, \vec{Y})\vec{Z} \cdot \vec{V} = (\nabla_{\vec{X}} \nabla_{\vec{Y}} - \nabla_{\vec{Y}} \nabla_{\vec{X}} - \nabla_{[\vec{X}, \vec{Y}]})\vec{Z} \cdot \vec{V} \quad (6)$$

by their "splitted representatives" (??), we have obtained the invariant generalizations of the Gauss-Codacci-Ricci's equations:

$$\begin{aligned} R(\vec{X}^\alpha, \vec{Y}^\alpha)\vec{Z}^\alpha \cdot \vec{V}^\alpha &= R^\alpha(\vec{X}^\alpha, \vec{Y}^\alpha)\vec{Z}^\alpha \cdot \vec{V}^\alpha + \\ &+ \sum_{\gamma \neq \alpha} \{2A^\gamma(\vec{X}^\alpha, \vec{Y}^\alpha) \cdot B^\gamma(\vec{Z}^\alpha, \vec{V}^\alpha) + \\ &+ B^\gamma(\vec{Y}^\alpha, \vec{V}^\alpha) \cdot B^\gamma(\vec{X}^\alpha, \vec{Z}^\alpha) - \\ &- B^\gamma(\vec{X}^\alpha, \vec{Y}^\alpha) \cdot B^\gamma(\vec{Y}^\alpha, \vec{Z}^\alpha)\}, \end{aligned} \quad (7)$$

where $R^\alpha(\vec{X}^\alpha, \vec{Y}^\alpha)\vec{Z}^\alpha \cdot \vec{V}^\alpha$ is the curvature tensor on the section Σ^α .

$$\begin{aligned} R(\vec{X}^\alpha, \vec{Y}^\alpha)\vec{Z}^\alpha \cdot \vec{V}^\beta &= \vec{V}^\beta \cdot \{(\nabla_{\vec{Y}^\alpha}^\beta \vec{B}^\beta)(\vec{X}^\alpha, \vec{Z}^\alpha) - \\ &- (\nabla_{\vec{X}^\alpha}^\beta \vec{B}^\beta)(\vec{Y}^\alpha, \vec{Z}^\alpha)\} - \\ &- \sum_{j \neq \alpha, \beta} Q^j(\vec{X}^\alpha, \vec{V}^\beta) \cdot B^j(\vec{Y}^\alpha, \vec{Z}^\alpha) + \\ &+ \sum_{j \neq \alpha, \beta} Q^j(\vec{Y}^\alpha, \vec{V}^\beta) \cdot B^j(\vec{X}^\alpha, \vec{Z}^\alpha) + \\ &+ \sum_{j \neq \alpha, \beta} 2\vec{Z}^\alpha \cdot Q^\alpha(\vec{A}^j(\vec{X}^\alpha, \vec{Y}^\alpha), \vec{V}^\beta) + \\ &+ 2\vec{Z}^\alpha \cdot B^\alpha(\vec{A}^\beta(\vec{X}^\alpha, \vec{Y}^\alpha), \vec{V}^\beta) \end{aligned} \quad (8)$$

$$\begin{aligned} R(\vec{X}^\alpha, \vec{Y}^\beta)\vec{Z}^\alpha \cdot \vec{V}^\beta &= (\vec{Z}^\alpha \cdot (\nabla_{\vec{X}^\alpha}^\alpha \vec{B}^\alpha) + \\ &+ (\vec{X}^\alpha \cdot \vec{B}^\alpha, \vec{Z}^\alpha \cdot \vec{B}^\alpha))(\vec{Y}^\beta, \vec{V}^\beta) + \\ &+ (\vec{V}^\beta \cdot (\nabla_{\vec{Y}^\beta}^\beta \vec{B}^\beta) + (\vec{Y}^\beta \cdot \vec{B}^\beta, \vec{V}^\beta \cdot \vec{B}^\beta))(\vec{X}^\alpha, \vec{Z}^\alpha) + \end{aligned}$$

$$\begin{aligned}
& + \sum_{\gamma \neq \alpha, \beta} \{B^\gamma(\vec{X}^\alpha, \vec{Z}^\alpha) \cdot B^\gamma(\vec{Y}^\beta, \vec{V}^\beta) - \\
& \quad - Q^\gamma(\vec{X}^\alpha, \vec{V}^\beta) \cdot Q^\gamma(\vec{Y}^\beta, \vec{Z}^\alpha) + \\
& \quad + \vec{V}^\beta \cdot Q^\beta(\vec{\Lambda}^\gamma(\vec{X}^\alpha, \vec{Y}^\beta), \vec{Z}^\alpha)\} \quad (9)
\end{aligned}$$

$$\begin{aligned}
R(\vec{X}^\alpha, \vec{Y}^\beta) \vec{Z}^\alpha \cdot \vec{V}^\delta = & \vec{V}^\delta \cdot \{(\nabla_{\vec{Y}^\beta}^\delta \vec{B}^\delta)(\vec{X}^\alpha, \vec{Z}^\alpha) - \\
& - (\nabla_{\vec{X}^\alpha}^\delta \vec{Q}^\delta)(\vec{Y}^\beta, \vec{Z}^\alpha)\} + \\
& + \vec{Z}^\alpha \cdot B^\alpha(\vec{Y}^\beta, \vec{Q}^\beta(\vec{X}^\alpha, \vec{V}^\delta)) - \\
& - \vec{V}^\delta \cdot B^\delta(\vec{Y}^\beta, \vec{B}^\beta(\vec{X}^\alpha, \vec{Z}^\alpha)) - \\
& - \vec{Z}^\alpha \cdot B^\alpha(\vec{\Lambda}^\alpha(\vec{X}^\alpha, \vec{Y}^\beta), \vec{V}^\delta) + \\
& + (\vec{Y}^\beta \cdot \vec{B}^\beta, \vec{V}^\delta \cdot \vec{B}^\delta)(\vec{X}^\alpha, \vec{Z}^\alpha) - \\
& - (\vec{X}^\alpha \cdot \vec{B}^\alpha, \vec{V}^\delta \cdot \vec{Q}^\delta)(\vec{Y}^\beta, \vec{Z}^\alpha) - \\
& - \sum_{\gamma \neq \alpha, \beta, \delta} \{\vec{Z}^\alpha \cdot Q^\alpha(\vec{\Lambda}^\gamma(\vec{X}^\alpha, \vec{Y}^\beta), \vec{V}^\delta) + \\
& + B^\gamma(\vec{X}^\alpha, \vec{Z}^\alpha) \cdot Q^\gamma(\vec{Y}^\beta, \vec{V}^\delta) + \\
& + Q^\gamma(\vec{X}^\alpha, \vec{V}^\delta) \cdot Q^\gamma(\vec{Y}^\beta, \vec{Z}^\alpha)\} \quad (10)
\end{aligned}$$

$$\begin{aligned}
R(\vec{X}^\alpha, \vec{Y}^\beta) \vec{Z}^\gamma \cdot \vec{V}^\delta = & \vec{V}^\delta \cdot \{(\nabla_{\vec{Y}^\beta}^\delta \vec{Q}^\delta)(\vec{X}^\alpha, \vec{Z}^\gamma) - \\
& - (\nabla_{\vec{X}^\alpha}^\delta \vec{Q}^\delta)(\vec{Y}^\beta, \vec{Z}^\gamma)\} + \vec{V}^\delta \{B^\delta(\vec{X}^\alpha, \vec{Q}^\alpha(\vec{Y}^\beta, \vec{Z}^\gamma)) - \\
& - B^\delta(\vec{Y}^\beta, \vec{Q}^\beta(\vec{X}^\alpha, \vec{Z}^\gamma)) - B^\delta(\vec{\Lambda}^\gamma(\vec{X}^\alpha, \vec{Y}^\beta), \vec{Z}^\gamma)\} + \\
& + \vec{Z}^\gamma \cdot \{B^\gamma(\vec{Y}^\beta, \vec{Q}^\beta(\vec{X}^\alpha, \vec{V}^\delta)) - \\
& - B^\gamma(\vec{X}^\alpha, \vec{Q}^\alpha(\vec{Y}^\beta, \vec{V}^\delta)) - B^\gamma(\vec{\Lambda}^\gamma(\vec{X}^\alpha, \vec{Y}^\beta), \vec{V}^\delta)\} + \\
& + (\vec{Y}^\beta \cdot \vec{B}^\beta, \vec{V}^\delta \cdot \vec{Q}^\delta)(\vec{X}^\alpha, \vec{Z}^\gamma) - \\
& - (\vec{X}^\alpha \cdot \vec{B}^\alpha, \vec{V}^\delta \cdot \vec{Q}^\delta)(\vec{Y}^\beta, \vec{Z}^\gamma) + \\
& + \sum_{\gamma \neq \alpha, \beta, \gamma, \delta} \{Q^j(\vec{Y}^\beta, \vec{V}^\delta) \cdot Q^j(\vec{X}^\alpha, \vec{Z}^\gamma) - \\
& - Q^j(\vec{Y}^\beta, \vec{Z}^\gamma) \cdot Q^j(\vec{X}^\alpha, \vec{V}^\delta) + \\
& + \vec{V}^\delta \cdot Q^\delta(\vec{\Lambda}^j(\vec{X}^\alpha, \vec{Y}^\beta), \vec{Z}^\gamma)\} \quad (11)
\end{aligned}$$

Here we use the following definitions:

$$\begin{aligned}
\Lambda^j(\vec{X}^\alpha, \vec{Y}^\beta) & \equiv H^j([\vec{X}^\alpha, \vec{Y}^\beta]) \\
(\vec{Y}^\beta \cdot \vec{B}^\beta, \vec{V}^\delta \cdot \vec{Q}^\delta)(\vec{X}^\alpha, \vec{Z}^\gamma) & \equiv \\
& \equiv g^{ab\alpha}[\vec{V}^\delta \cdot Q^\delta(\vec{Z}^\gamma, \vec{e}_a^\alpha) \vec{X}^\beta] \cdot B^\beta(\vec{X}^\alpha, \vec{e}_b^\alpha) \quad (12) \\
(\nabla_{\vec{X}^\beta}^\alpha \vec{Q}^\alpha)(\vec{Y}^\gamma, \vec{V}^\delta) & \equiv \nabla_{\vec{X}^\beta}^\alpha \vec{Q}^\alpha(\vec{Y}^\gamma, \vec{V}^\delta) - \\
& - \vec{Q}^\alpha(\nabla_{\vec{X}^\beta}^\gamma \vec{Y}^\gamma, \vec{V}^\delta) - \vec{Q}^\alpha(\vec{Y}^\alpha, \nabla_{\vec{X}^\beta}^\delta \vec{V}^\delta)
\end{aligned}$$

Objects S and A are the symmetric and antisymmetric parts of B respectively. All popular splitting methods (such as the monad method Antonov V.I., et.al., 1978, the Newman-Penrose's formalism Newman E.T. and Penrose R., 1962, the Geroch's formalism Geroch R., 1971) and many others can be obtained as a special cases of this approach when we choose concrete basis of splitting.

For example, the Newman-Penrose's method corresponds to the following projection operators:

$$\begin{aligned}
H' = E_a \otimes \theta^a, \quad H'' = I - E_a \otimes \theta^a, \\
g = g' + g'' \quad (13)
\end{aligned}$$

$$\begin{aligned}
g' = 2(\theta^{\{1} \otimes \theta^{\{2} + \theta^{\{3} \otimes \theta^{\{3} + \dots \\
+ \theta^{\{n-1} \otimes \theta^{\{n}\})} \quad (14)
\end{aligned}$$

The basis of one-forms $\{\theta^a\}$, $a = 1, 2, \dots, n$ is a complex one in general case.

References

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