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Multifunctional polarization tomography of birefringence networks of biological crystals

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Abstract. The results of optical modeling of biological crystals polycrystalline multilayer networks have been presented. Determined in this work are the algorithms of reconstruction of parameter distributions that describe the linear and circular birefringence. To separate manifestations of these mechanisms, we propose the method of space-frequency filtering. Criteria for differentiation of benign and malignant tissues of the women reproductive sphere have been found.

Keywords: polarization tomography, biological crystals, diagnostics.

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1. Introduction

In numerous areas of optical diagnostics of biological tissue structures, the techniques based on polarized laser radiation are widely used [1-3]. Particularly, laser polarimetry of microscopic images of polycrystalline protein networks was formed as a separate approach in the study of optically anisotropic component of different biological tissues biopsy [4-28].

Our research is aimed at designing the experimental method of Fourier polarimetry and spatial-frequency selection of parameter distributions of linear and circular birefringence of the biopsy of the uterine wall tissue in order to differentiate benign (fibromyoma) and malignant (adenocarcinoma) conditions.

2. Theory of the method

The following model concepts have been assumed [14] by us as the basis for an analysis of the processes of amplitude-phase modulation of laser radiation by the polycrystalline network of biological crystals with linear and circular birefringence:

• the uterine wall tissue consists of two optically anisotropic layers of the muscular (myometrium) and connective (endometrium) layers;

- phase anisotropy of these layers formed by optically active myosin and collagen proteins that form fibrillar network with different conformational structure;
- the optically anisotropic myometrium network is formed by spatially ordered myosin fibrils. This spatial organization generates large-scale (the range of the transverse dimensions $l \sim 50...300 \, \mu \text{m}$) optically uniaxial birefringent linear structure. Polarization properties of biological crystals are characterized by the Jones matrix of the following form

$$\{Q\} = \left\| \frac{\sin^2 \rho + \cos^2 \rho \exp(-i\delta)}{\sin \rho \cos \rho (1 - \exp(-i\delta))} \right\|,$$

$$\sin \rho \cos \rho (1 - \exp(-i\delta)); \quad \cos^2 \rho + \sin^2 \rho \exp(-i\delta) \right\|,$$
(1)

where ρ is the direction of the axes, it determines the direction of stacking fibrils; $\delta = \frac{2\pi l}{\lambda} \Delta n$ is the value of the phase shift between orthogonal components of the amplitude; λ – wavelength; Δn – linear birefringence;

 the optically anisotropic component of the endometrium is formed by small-scale (l ~ 5...25 μm) randomly oriented collagen fibrils with predominant circular birefringence

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$$\{A\} = \begin{vmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{vmatrix}, \tag{2}$$

where $\theta = \frac{\pi l}{\lambda} \Delta n_{\pm}$ is the angle of the polarization plane;

 Δn_+ – indicator of the circular birefringence.

• the process of converting laser radiation by optically anisotropic fibrils is characterized by the following matrix equation

$$E = \{Q\}\{A\}E_0. \tag{3}$$

The coordinate distribution of parameters describing linear δ and circular θ anisotropy of fibrillar network layers of the myometrium and endometrium can be found from the corresponding Jones matrix (relations (4) and (5)).

$$\delta(q_{ik}) = \arccos \left\{ \frac{i \left(E_{y} E_{x}^{*} - E_{x} E_{y}^{*} \right)_{\otimes}}{\left(E_{y} E_{x}^{*} + E_{x} E_{y}^{*} \right)_{\otimes}} - \right.$$

$$\left. - 0.5i \left[\frac{\left(E_{y} E_{x}^{*} - E_{x} E_{y}^{*} \right)_{0}}{\left(E_{y} E_{x}^{*} - E_{x} E_{y}^{*} \right)_{0}} + \frac{\left(E_{y} E_{x}^{*} - E_{x} E_{y}^{*} \right)_{90}}{\left(E_{y} E_{x}^{*} + E_{x} E_{y}^{*} \right)_{90}} \right] \right\},$$

$$(4)$$

$$\theta(a_{ik}) = \arccos \left\{ \frac{\left(E_x E_x^* - E_y E_y^*\right)_0}{\left(E_x E_x^* + E_y E_y^*\right)_0} - \frac{\left(E_x E_x^* - E_y E_y^*\right)_{90}}{\left(E_x E_x^* + E_y E_y^*\right)_{90}} \right\}. (5)$$

Here, we use the following notations

- * complex conjugate;
- \otimes sensing the right circularly polarized wave $E_0 = \begin{pmatrix} 1 \\ i \end{pmatrix}$;
- 0 sounding the linearly polarized wave $E_0 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ with azimuth 0° ;
- 90 sounding the linearly polarized wave $E_0 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ with azimuth 90°.

From the medical point of view, "optic separation" of various anisotropy mechanisms in layers of the uterine wall is the topical task. The fact is that at the initial, early stages, the pathological changes are accompanied by a concentration increase in the amount of optically active proteins in polycrystalline network of the myometrium and endometrium. Fibrillar structural changes – the formation of the growth direction of tumour shoots and so on – correspond to a more developed pathological condition.

With the purpose to select these conditions, we have applied the method of spatial-frequency filtration of the polarization-inhomogeneous field $E_x(\rho, \delta, \theta)$ and $E_y(\rho, \delta, \theta)$ in Fourier plane. The main idea of this approach is that spatial-frequency structure of Fourier form of laser image inherent to the tissue of the uterine wall is different for its large-scale (myosin) and small-scale (collagen) protein structures. Proceeding from this,

using spatial-frequency filtration, it is mainly possible to isolate either low-frequency (with linear birefringence) or high-frequency (with circular birefringence) components, which by means of Fourier conversion can be transformed into respectively "separated" laser images.

In our case, we consider the situation when the distance between projection lens, the layer of the object under study, and the plane of registration constitute the focal distance f.

$$U_x(v,\mu) = \frac{1}{i\lambda f} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} E_x(x,y) \exp[-i2\pi(xv + y\mu)] dxdy, \quad (6)$$

$$U_{y}(v,\mu) = \frac{1}{i\lambda f} \int_{-\infty}^{\infty} \int E_{y}(x,y) \exp[-i2\pi(xv + y\mu)] dxdy.$$
 (7)

Here, U_x , U_y are Fourier images of distributions $E_x(\rho, \delta, \theta)$ and $E_y(\rho, \delta, \theta)$ in the focal plane with spatial

frequencies
$$v = \frac{x}{\lambda f}$$
 and $\mu = \frac{y}{\lambda f}$.

If we place vignetting transparent $R(\Delta v, \Delta \mu)$ or opaque $R^{-1}(\Delta v, \Delta \mu)$ diaphragm in the central part of Fourier plane, we can identify corresponding areas of the spatial-frequency structure of the distribution amplitudes, which are formed by predominant influences of linear $\hat{U}(\rho, \delta, v, \mu)$ and circular $\hat{U}(\theta, v, \mu)$ birefringence

$$\begin{cases}
\hat{U}(\rho, \delta, \nu, \mu) = R(\Delta \nu, \Delta \mu) U(\nu, \mu) \\
\hat{U}(\theta, \nu, \mu) = R^{-1}(\Delta \nu, \Delta \mu) U(\nu, \mu).
\end{cases}$$
(8)

To renew the proper distributions of complex

amplitudes
$$\begin{cases} \hat{E}_x(\rho, \delta, x, y) \\ \dot{E}_x(\theta, x, y) \end{cases} \text{ and } \begin{cases} \hat{E}_y(\rho, \delta, x, y) \\ \dot{E}_y(\theta, x, y) \end{cases} \text{ in the }$$

image plane of this layer is possible by reverse Fourier transformation

$$\begin{cases}
\hat{E}_{x}(\rho,\delta,x,y) = \\
= \frac{1}{i\lambda f} \int_{-\infty}^{\infty} \int R(\Delta v, \Delta \mu) \hat{U}_{x}(v,\mu) \exp\left[i2\pi(x\Delta v + y\Delta \mu)\right] dv d\mu, \\
\hat{E}_{x}(\theta,x,y) = \\
= \frac{1}{i\lambda f} \int_{-\infty}^{\infty} \int R^{-1}(\Delta v, \Delta \mu) \hat{U}_{x}(v,\mu) \exp\left[i2\pi(x\Delta v + y\Delta \mu)\right] dv d\mu, \\
\hat{E}_{y}(\rho,\delta,x,y) = \\
= \frac{1}{i\lambda f} \int_{-\infty}^{\infty} \int R(\Delta v, \Delta \mu) \hat{U}_{y}(v,\mu) \exp\left[i2\pi(x\Delta v + y\Delta \mu)\right] dv d\mu, \\
\hat{E}_{y}(\theta,x,y) = \\
= \frac{1}{i\lambda f} \int_{-\infty}^{\infty} \int R^{-1}(\Delta v, \Delta \mu) \hat{U}_{y}(v,\mu) \exp\left[i2\pi(x\Delta v + y\Delta \mu)\right] dv d\mu. \\
\end{cases}$$
(10)

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Within each pixel of a digital camera kits measured intensities $R \to I_{\otimes}^{\otimes}$; I_{\oplus}^{\otimes} ; I_{\otimes}^{0} ; I_{\oplus}^{0} ; I_{\otimes}^{90} ; I_{\oplus}^{90} and $R^{-1} \to I_{0}^{0}$; I_{90}^{0} ; I_{0}^{90} ; I_{90}^{90} . The magnitude of phase shift $\delta(R)$ and a rotation of plane polarization $\theta(R^{-1})$ are calculated according to the following equations.

$$\delta(R) = \arccos \left\{ \frac{\left(I_{\otimes}^{\otimes} - I_{\oplus}^{\otimes}\right)}{\left(I_{\otimes}^{\otimes} + I_{\oplus}^{\otimes}\right)} - \\ -0.5 \left[\frac{\left(I_{\otimes}^{0} - I_{\oplus}^{0}\right)}{\left(I_{\otimes}^{0} + I_{\oplus}^{0}\right)} + \frac{\left(I_{\otimes}^{90} - I_{\oplus}^{90}\right)}{\left(I_{\otimes}^{90} + I_{\oplus}^{90}\right)} \right] \right\},$$

$$(11)$$

$$\theta\left(R^{-1}\right) = 0.5 \arccos\left\{\frac{\left(I_0^0 - I_{90}^0\right)}{\left(I_0^0 + I_{90}^0\right)} - \frac{\left(I_0^{90} - I_{90}^{90}\right)}{\left(I_0^{90} + I_{90}^{90}\right)}\right\}. \tag{12}$$

To estimate distributions $q = \begin{cases} \delta(m, n) \\ \theta(m, n) \end{cases}$, we calculated a set of statistical moments of the 1st to 4th orders $Z_{i=1:2:3:4}$ using the following algorithms [2]

$$Z_{1} = \frac{1}{N} \sum_{i=1}^{N} |(q)_{i}|, \ Z_{2} = \sqrt{\frac{1}{N} \sum_{i=1}^{N} (q)_{i}^{2}},$$

$$Z_{3} = \frac{1}{(Z_{2})^{3}} \frac{1}{N} \sum_{i=1}^{N} (q)_{i}^{3}, \ Z_{4} = \frac{1}{(Z_{2})^{4}} \frac{1}{N} \sum_{i=1}^{N} (q)_{i}^{4}.$$
(13)

3. Experimental results and discussion

As objects of investigation, we chose two groups of optically-thin (extinction coefficient $\tau \approx 0.090...0.095$) histological sections of the biopsy taken from benign (26 samples) and malignant (26 samples) tumor of the uterus. The samples were prepared on a freezing microtome by using the standard technique.

3.1. Spatial frequency Fourier polarimetry of linear birefringence of the fibrillar matrices of the uterine tissue

For the purpose of choosing optimal conditions of spatial-frequency filtration, the following range $\Delta r = 2...50$ pix of possible dimensions of the vignetting diaphragm. As a criterion for choosing the size of the aperture was both an extreme change in the set of statistical moments of the 1st to 4th order $Z_{j=1;2;3;4}$ (relation (13)) that characterize the distribution $\delta(v,\mu)$. In our case, we obtained the following optimal size $\Delta r = 30$ pix. This geometric size was chosen for comparative investigations of optical anisotropy of the fibrillar networks of the uterine tissue with benign and malignant changes (Fig. 1).

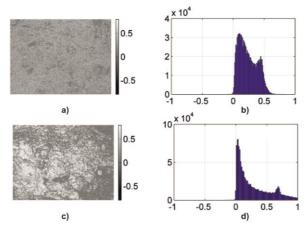


Fig. 1. "Low-frequency" coordinate maps (a, c) and histograms (b, d) of the linear birefringence distribution of a histological section of benign uterine tumor (fibromyoma) biopsy (a-b) and malignant uterine tumor (adenocarcinoma) biopsy (c-d).

Quantitative differences between phase maps of the protein matrices of the uterine tissue of both types illustrate average values and standard deviations of the set of statistic $Z_{i=1;2;3;4}$ parameters, listed in Table 1. Separate figures are given specificity $Sp = a (a + b)^{-1} (a, b - \text{numbers of correct and incorrect conclusions)}$ of these parameters in differentiation of benign and malignant uterine tumors.

A comparative analysis of the data of Fourier polarimetry of linear birefringence of protein polycrystalline networks of histological section taken from benign and malignant uterine tumors biopsy revealed sensibility to differentiation of these states $(Z_{i=1;2;3;4})$ with the specificity level $Sp \sim 63\%...79\%$ (printed in grey color in Table 1).

3.2. Special frequency Fourier polarimetry of circular birefringence of the fibrillar networks of uterine tissue

Diagnostic possibilities of differentiating the pathologically changed samples of uterine tissue illustrate probable adduced in a series in Fig. 2.

Quantitative differences between phase maps of circular birefringence of the endometrium collagenous networks illustrate average values and standard deviations of the statistic set $Z_{i=1,2;3;4}$ parameters, adduced in Table 2.

Table 1. Parameters of the statistical, correlative and selfsimilar structure of phase maps for linear birefringence of uterine tissue networks.

Parameters	$\delta(m \times n)$		Cn
rarameters	Fibromyoma	Adenocarcinoma	Sp
Z_1	0.37±0.059	0.47±0.11	71%
Z_2	0.16±0.022	0.19±0.031	63%
Z_3	1.16±0.37	0.73±0.14	77%
Z_4	0.79±0.21	0.57±0.051	79%

Table 2. Parameters of the statistical, correlative and selfsimilar structure of phase maps for circular birefringence of polycrystalline uterine tissue networks.

Parameters	$\Theta(m \times n)$		Sp
Farameters	Fibromyoma	Adenocarcinoma	sp
Z_1	0.11±0.021	0.21±0.037	73%
Z_2	0.14±0.018	0.23±0.035	81%
Z_3	0.79±0.11	1.38±0.39	86%
Z_4	0.81±0.13	1.21±0.21	83%

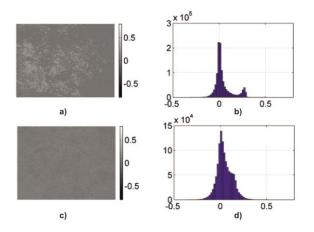


Fig. 2. "High-frequency" coordinate maps (a, c) and histograms (b, d) of the linear birefringence distribution of a histological section of benign uterine tumor (fibromyoma) biopsy (a-b) and malignant uterine tumor (adenocarcinoma) biopsy (c-d).

A comparative analysis revealed the following oncologic condition with a high level of specificity $(73\% \le Sp \le 86\%)$ of uterus parameters (depicted in grey color in Table 2):

• statistical moments of the 1st and 2nd orders $\theta(m \times n)$, differences between the values lies within the range 2 to 3.5 times.

4. Conclusions

The model of generalized optic anisotropy in the fibrillar protein matrices has been suggested, and the method of Fourier polarimetry for the parameters of linear and circular birefringence of the tissue biopsy of the uterine wall has been substantiated.

The set of criteria for differentiation of benign (fibromyoma) and malignant (adenocarcinoma) conditions of the uterine has been revealed and substantiated:

- statistical moments of the 1st–4th orders that characterize distribution of the phase shift between orthogonal components of the laser radiation are defined by linear birefringence of myosin fibers of the endometrium layer;
- statistical moments of the 1^{st} 4^{th} orders that characterize distribution of plane rotations of laser

radiation polarization are defined by circular birefringence of the network of collagenic fibrils of the endometrium layer.

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