



SOME PROPERTIES OF BOUNDARY VALUE PROBLEMS FOR BESSEL'S EQUATION

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Some properties of systems of the Bessel functions of negative order less than -1 generated by one boundary value problem are studied.

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Досліджено деякі властивості систем функцій Бесселя з від'ємним індексом меншим за -1 , які виникають в одній краївій задачі.

Let $\nu > 0$ be a non-integer number, $\overline{0;\nu} = \mathbb{Z} \cap [0; \nu]$, $p \in \mathbb{R}$, and $L_2((0; 1); x^p dx)$ be the Hilbert space of functions $f : (0; 1) \rightarrow \mathbb{C}$ such that the function $t^{p/2}f(t)$ belongs to the space $L_2(0; 1)$; the inner product and the norm in $L_2((0; 1); x^p dx)$ are given respectively by $\langle f_1; f_2 \rangle = \int_0^1 t^p f_1(t) \overline{f_2(t)} dt$ and $\|f\| = \sqrt{\int_0^1 t^p |f(t)|^2 dt}$. Further, let J_ν be Bessel's function of order ν . We study the approximation properties of the system $\{J_{-\nu}(\rho_j x)\}$ for some sequence (ρ_j) tending to infinity. Such a system arises, for example, when considering the boundary value problem

$$-f'' + \frac{\nu^2 - 1/4}{x^2} f = \lambda f, \quad (1)$$

$$f(1) = 0, \quad (2)$$

$$\exists c_k \in \mathbb{C} : \quad f(x) = \sum_{k \in \overline{0;\nu}} c_k x^{-\nu+2k+1/2} + o(x^{\nu+1/2}), \quad x \rightarrow 0+. \quad (3)$$

Indeed, one easily derives the following result:

Proposition 1. *The boundary value problem (1)–(3) has a countable set of eigenvalues $\{\lambda_j : j \in \mathbb{N}\}$, and $\lambda_j = \rho_j^2$, where ρ_j are the zeros of the Bessel function $J_{-\nu}$; moreover, $v_j(x) = \rho_j^\nu \sqrt{\pi x/2} J_{-\nu}(\rho_j x)$ are the corresponding eigenfunctions.*

Proof. For $\lambda \neq 0$, equation (1) has two linearly independent solutions given by

$$u(x) = \rho^{-\nu-1/2} \sqrt{\pi \rho x/2} J_\nu(\rho x), \quad v(x) = \chi_\nu \rho^{\nu-1/2} \sqrt{\pi \rho x/2} J_{-\nu}(\rho x),$$

where $\rho = \sqrt{\lambda}$ and $\chi_\nu = -1/\sin(\nu\pi)$. Using the asymptotic behaviour of the Bessel functions, we find that

$$\begin{aligned} u(x) &= \sqrt{\pi/2} \frac{x^{\nu+1/2}}{2^\nu \Gamma(\nu+1)} + o(x^{\nu+1/2}), \quad x \rightarrow 0+, \\ v(x) &= \chi_\nu \sqrt{\pi/2} \sum_{k \in \overline{0}; \nu} \frac{(-1)^k \rho^{2k} x^{-\nu+2k+1/2}}{2^{-\nu+2k} k! \Gamma(-\nu+k+1)} + O(x^{-\nu+2[\nu]+5/2}), \quad x \rightarrow 0+. \end{aligned}$$

Therefore the general solution $\alpha u(x) + \beta v(x)$ of (1) satisfies the condition (3) if and only if $\alpha = 0$. Recalling (2), we conclude that the eigenfunctions of the boundary value problem (1)–(3) are $v_j(x) = \rho_j^\nu \sqrt{\pi x/2} J_{-\nu}(\rho_j x)$, and the corresponding eigenvalues are $\{\lambda_j : j \in \mathbb{N}\}$, where $\lambda_j = \rho_j^2$ and ρ_j are the zeros of the Bessel function $J_{-\nu}$. The proposition is proved. \square

The function $J_{-3/2}$ has ([1]) an infinite set $\{\rho_j : j \in \mathbb{Z} \setminus \{0\}\}$ of zeros; moreover, the zeros ρ_1 and $\rho_{-1} = \bar{\rho}_1 = -\rho_1$ are purely imaginary, ρ_j with $j \in \mathbb{N} \setminus \{1\}$ are positive, and $\rho_{-j} := -\rho_j$, $j \in \mathbb{N} \setminus \{1\}$, are negative. In [2]–[4] the following results were proved.

Theorem A. *Let ρ_j be the zeros of the function $J_{-3/2}$. Then the system $\{\rho_j \sqrt{\rho_j x} J_{-3/2}(\rho_j x) : j \in \mathbb{N} \setminus \{1\}\}$ is complete in the space $L_2((0; 1); x^2 dx)$, has in this space a biorthogonal system $\{\gamma_k : k \in \mathbb{N} \setminus \{1\}\}$,*

$$\overline{\gamma_k}(x) := \frac{\pi(\rho_k \sqrt{\rho_k x} J_{-3/2}(\rho_k x) - \rho_1 \sqrt{\rho_1 x} J_{-3/2}(\rho_1 x))}{x^2 \rho_k^2 \cos^2 \rho_k},$$

and is not a basis of this space.

Theorem B. *For $\nu = 3/2$ the boundary value problem (1)–(3) has a countable set of eigenvalues $\{\lambda_j : j \in \mathbb{N}\}$; moreover, $\lambda_j = \rho_j^2$, where ρ_j are the zeros of the Bessel function $J_{-3/2}$. In particular, all eigenvalues are real, λ_1 is negative, and λ_j are positive for $j > 1$. The corresponding eigenfunctions $\rho_j \sqrt{\rho_j x} J_{-3/2}(\rho_j x)$ enjoy the properties described in Theorem A.*

Similar results can be obtained for $\nu = 5/2$. Their extension to arbitrary ν leads to some difficulties. In this connection we will prove here the following proposition.

Theorem 1. *Let $z \in \mathbb{C}$ and let ρ_j be the zeros of the function $J_{-\nu}$ with $\nu = m + 1/2$, $m \in \mathbb{N}$. Then*

$$\int_0^1 \cos(zt) \frac{J_{-\nu}(\rho_j t)}{(\rho_j t)^{-\nu}} dt = (-1)^{m+1} \sqrt{\frac{\pi}{2}} \rho_j^{1+\nu} J_{-\nu+1}(\rho_j) \left(\frac{1}{z} \frac{d}{dz} \right)^m \frac{J_{-\nu}(z)}{(z^2 - \rho_j^2) z^{-\nu}},$$

where $(z^{-1} d/dz)^m$ is the m -th power of the differential operator $z^{-1} d/dz$.

Proof. The functions $f(t) = J_{\pm\nu}(zt)$ are the solutions of the equation $f'' + f'/t - \nu^2 f/t^2 = -z^2 f$. Hence

$$\begin{aligned} \frac{d}{dt} \left(t \frac{dJ_{-\nu}(zt)}{dt} \right) + \left(z^2 t - \frac{\nu^2}{t} \right) J_{-\nu}(zt) &= 0, \\ \frac{d}{dt} \left(t \frac{dJ_{-\nu}(\rho_j t)}{dt} \right) + \left(\rho_j^2 t - \frac{\nu^2}{t} \right) J_{-\nu}(\rho_j t) &= 0. \end{aligned}$$

Multiply the first equality by $J_{-\nu}(\rho_j t)$, the second by $J_{-\nu}(zt)$ and subtract. Then

$$\frac{d}{dt} \left(t \left(\frac{dJ_{-\nu}(zt)}{dt} J_{-\nu}(\rho_j t) - \frac{dJ_{-\nu}(\rho_j t)}{dt} J_{-\nu}(zt) \right) \right) = -(z^2 - \rho_j^2) t J_{-\nu}(zt) J_{-\nu}(\rho_j t)$$

or, in turn,

$$\rho_j^{-\nu} t^{-2\nu+1} \frac{J_{-\nu}(zt)}{(zt)^{-\nu}} \frac{J_{-\nu}(\rho_j t)}{(\rho_j t)^{-\nu}} = \frac{d}{dt} \left(t \frac{-(J_{-\nu}(zt))'_t J_{-\nu}(\rho_j t) + (J_{-\nu}(\rho_j t))'_t J_{-\nu}(zt)}{(z^2 - \rho_j^2) z^{-\nu}} \right),$$

and

$$\begin{aligned} \rho_j^{-\nu} t^{-2\nu+1} \frac{J_{-\nu}(\rho_j t)}{(\rho_j t)^{-\nu}} \left(\frac{1}{z} \frac{\partial}{\partial z} \right)^m \frac{J_{-\nu}(zt)}{(zt)^{-\nu}} &= \\ &= \left(\frac{1}{z} \frac{\partial}{\partial z} \right)^m \frac{\partial}{\partial t} \left(t \frac{-(J_{-\nu}(zt))'_t J_{-\nu}(\rho_j t) + (J_{-\nu}(\rho_j t))'_t J_{-\nu}(zt)}{(z^2 - \rho_j^2) z^{-\nu}} \right). \end{aligned}$$

Using the equality [5, p.56]

$$\frac{d}{dz} \frac{J_{-\nu}(z)}{z^{-\nu}} = -z^\nu J_{-\nu+1}(z),$$

we obtain

$$\left(\frac{1}{z} \frac{\partial}{\partial z} \right)^m \frac{J_{-\nu}(zt)}{(zt)^{-\nu}} = (-1)^m t^{2m} \frac{J_{-\nu+m}(zt)}{(zt)^{-\nu+m}}.$$

Therefore

$$\begin{aligned} (-1)^m \rho_j^{-\nu} \frac{J_{-\nu}(\rho_j t)}{(\rho_j t)^{-\nu}} t^{-2\nu+1+2m} \frac{J_{-\nu+m}(zt)}{(zt)^{-\nu+m}} &= \\ &= \left(\frac{1}{z} \frac{\partial}{\partial z} \right)^m \frac{\partial}{\partial t} \left(t \frac{-(J_{-\nu}(zt))'_t J_{-\nu}(\rho_j t) + (J_{-\nu}(\rho_j t))'_t J_{-\nu}(zt)}{(z^2 - \rho_j^2) z^{-\nu}} \right). \end{aligned}$$

Since $-\nu + m = -1/2$ and $t^{1/2} J_{-1/2}(t) = \sqrt{2/\pi} \cos t$, we have

$$\begin{aligned} (-1)^m \rho_j^{-\nu} \sqrt{\frac{2}{\pi}} \frac{J_{-\nu}(\rho_j t)}{(\rho_j t)^{-\nu}} \cos(zt) &= \\ &= \left(\frac{1}{z} \frac{\partial}{\partial z} \right)^m \frac{\partial}{\partial t} \left(t \frac{-(J_{-\nu}(zt))'_t J_{-\nu}(\rho_j t) + (J_{-\nu}(\rho_j t))'_t J_{-\nu}(zt)}{(z^2 - \rho_j^2) z^{-\nu}} \right) = \\ &= \frac{\partial}{\partial t} \left(\frac{1}{z} \frac{\partial}{\partial z} \right)^m \left(t \frac{-(J_{-\nu}(zt))'_t J_{-\nu}(\rho_j t) + (J_{-\nu}(\rho_j t))'_t J_{-\nu}(zt)}{(z^2 - \rho_j^2) z^{-\nu}} \right). \end{aligned}$$

Hence

$$\begin{aligned}
 (-1)^m \rho_j^{-\nu} \sqrt{\frac{2}{\pi}} \int_0^1 \cos(zt) \frac{J_{-\nu}(\rho_j t)}{(\rho_j t)^{-\nu}} dt &= \\
 &= \int_0^1 \frac{\partial}{\partial t} \left(\frac{1}{z} \frac{\partial}{\partial z} \right)^m \left(t \frac{-(J_{-\nu}(zt))'_t J_{-\nu}(\rho_j t) + (J_{-\nu}(\rho_j t))'_t J_{-\nu}(zt)}{(z^2 - \rho_j^2) z^{-\nu}} \right) dt = \\
 &= \left(\frac{1}{z} \frac{\partial}{\partial z} \right)^m t \frac{-(J_{-\nu}(zt))'_t J_{-\nu}(\rho_j t) + (J_{-\nu}(\rho_j t))'_t J_{-\nu}(zt)}{(z^2 - \rho_j^2) z^{-\nu}} \Big|_0^1 = \\
 &= \rho_j J'_{-\nu}(\rho_j) \left(\frac{1}{z} \frac{d}{dz} \right)^m \frac{J_{-\nu}(z)}{(z^2 - \rho_j^2) z^{-\nu}}.
 \end{aligned}$$

But ([5, p.56]) $t J'_{-\nu}(t) + \nu J_{-\nu}(t) = -t J_{-\nu+1}(t)$. Then $J'_{-\nu}(\rho_j) = -J_{-\nu+1}(\rho_j)$, and we obtain the desired result. The theorem is proved. \square

Theorem 1 for $\nu = 3/2$ was used to prove completeness of the corresponding systems of eigenfunction in Theorem B.

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