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EFFECTS OF PROXIMITY POTENTIALS ON THE CROSS-SECTIONS OF ^{6,8}He + ^{65}Cu HALO FUSION REACTIONS

UDC 539

The comprehensive theoretical study is performed to determine the best proximity potentials in reproducing ^{6,8}He + ^{65}Cu fusion reactions. Twenty three different versions of proximity potentials that consist of Prox 66, Prox 76, Prox 77, Prox 79, Prox 81-I, Prox 81-II, Prox 81-III, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003-I, Prox 2003-II, Prox 2003-III, Prox 2010, BW 91, AW 95, Bass 73, Bass 77, Bass 80, CW 76, Ngo 80, and D are used. The theoretical results are compared with experimental data on ^{6,8}He + ^{65}Cu fusion reactions. The appropriate proximity potentials are determined.

K e y w o r d s: fusion cross-sections, proximity potentials, halo nuclei.

1. Introduction

Studies of halo nuclei are still a current topic in the field of nuclear physics. It is well known that the parameters such as the binding energy, isospin values, wave function(s) of the valence nucleon(s), and density distributions are different for the halo nuclei; ^{6}He, ^{8}He, ^{11}Li, ^{11}Be, ^{14}Be, ^{14}B, ^{15}C, and ^{19}C are accepted to be neutron halo nuclei; ^{6}He and ^{8}He are among the most important of these nuclei. In addition to the elastic scattering interactions of these nuclei, the fusion reactions are also an important issue. In this context, Navin et al. [1] have presented the data on ^{6}He + ^{65}Cu fusion reaction. Chatterjee *et al.* [2] the displayed exclusive measurement of fusion ^{6}He by ^{65}Cu. Then the fusion cross-section of ^{8}He + ^{65}Cu reaction has been reported by Lemasson *et al.* [3]. Both theoretical and experimental studies for understanding the fusion processes of halo nuclei ^{6}He [4–7] and ^{8}He [8, 9] can be found in the literature.

The nucleus-nucleus potentials are very important to study the fusion reactions of halo nuclei. For this purpose, different potentials such as those in the double folding model [10], time-dependent Hartree–Fock theory [11], and model with Skyrme energy-density functional [12] can be evaluated. In addition, proximity potentials are of importance in obtaining the nuclear potentials of fusion reactions. The alpha decay process [13], cluster decay process [14], fission reaction [15], and elastic scattering reactions [16–19] have been examined, by using these potentials. On the other hand, it can be seen, while the literature is examined, that the fusion reactions of ^{6}He and ^{8}He nuclei have not been analyzed, by using a wide range of proximity potentials. Therefore, we consider that it will be useful and interesting to overcome this deficiency in the literature for ^{6}He and ^{8}He halo nuclei which are isotopes with each other.

In the present study, we will examine the effects of the proximity potentials on ^{6}He + ^{65}Cu and ^{8}He + ^{65}Cu fusion cross-sections. It would be interesting to compare the fusion results of two halo nuclei with the neutron rich isotopes over the same target nucleus. With this goal, we apply twenty three different types of proximity potentials to obtain the fusion cross-sections, which can be emphasized as proximity 1966, proximity 1976, proximity 1977, proximity

1979, proximity 1981-I, proximity 1981-II, proximity 1981-III, proximity 1984, proximity 1988, modified proximity 1988, proximity 1995, proximity 2003-I, proximity 2003-II, proximity 2003-III, proximity 2010, Broglia and Winther 1991, Aage Winther, and Bass 1973, Bass 1977, Bass 1980, Christensen and Winther 1976, Ngo 1980, and Denisov potential. Then we compare the theoretical results and experimental data.

Section 2 gives a brief description of proximity potentials used in the theoretical calculations. Section 3 presents the calculation procedure. Section 4 shows the results and the discussion of calculations. Section 5 provides the summary and conclusion.

2. Proximity Potentials

In the theoretical analysis of fusion cross sections of ^{6,8}He + ^{65}Cu reactions, we apply twenty three different versions of proximity potentials. These potentials are explained in the following subsections.

2.1. Prox 66, Prox 76, Prox 77, Prox 79, Prox 81, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003, Prox 2010 potentials

Prox 77 version of the proximity potential [20, 21] is in the following form:

$$V_{\rm N}^{\rm Prox\,77}(r) = 4\pi\gamma b\overline{R}\Phi\left(\zeta = \frac{r - C_1 - C_2}{b}\right) \text{ MeV}, \quad (1)$$

where

$$\overline{R} = \frac{C_1 C_2}{C_1 + C_2}, \;\; C_i = R_i\left[1 - \left(\frac{b}{R_i}\right)^2 + ...\right], \;\; b \approx 1\,\text{fm}, \quad (2)$$

R_i (effective radius) has the form

$$R_i = 1.28A_i^{1/3} - 0.76 + 0.8A_i^{-1/3}\,\text{fm} \quad (i = 1, 2), \quad (3)$$

and γ (surface energy coefficient) is taken as

$$\gamma = \gamma_0\left[1 - k_s\left(\frac{N-Z}{N+Z}\right)^2\right], \quad (4)$$

where $N(Z)$ is the total number of neutrons(protons). The universal function $\Phi(\zeta)$ is evaluated as

$$\Phi(\zeta) = \begin{cases} -\frac{1}{2}(\zeta - 2.54)^2 - 0.0852(\zeta - 2.54)^3 & \text{for } \zeta \le 1.2511, \\ -3.437\exp\left(-\frac{\zeta}{0.75}\right) & \text{for } \zeta \ge 1.2511. \end{cases}$$

Table 1. **γ_0 and k_s values of Prox 66, Prox 76, Prox 77, Prox 79, Prox 81, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003, and Prox 2010 potentials**

Potential type	γ_0, MeV/fm^2	k_s	Ref.
Prox 66	1.01734	1.79	[22]
Prox 76	1.460734	4.0	[23]
Prox 77	0.9517	1.7826	[22]
Prox 79	1.2402	3.0	[24]
Prox 81-I	1.1756	2.2	[25]
Prox 81-II	1.27326	2.5	[25]
Prox 81-III	1.2502	2.4	[25]
Prox 84	0.9517	2.6	[26]
Prox 88	1.2496	2.3	[27]
Mod-Prox-88	1.65	2.3	[28]
Prox 95	1.25284	2.345	[29]
Prox 2003-I	1.08948	1.9830	[30]
Prox 2003-II	0.9180	0.7546	[30]
Prox 2003-III	0.911445	2.2938	[30]
Prox 2010	1.460734	4.0	[21, 31, 32]

Then a lot of studies have been performed on the proximity potentials. Different values of γ_0 and k_s have been proposed as a result of these studies. The other parameters of these potentials are the same as for Prox 77 potential. In this context, there are fifteen various potentials examined in this work. The γ_0 and k_s values of these potentials values are given in Table 1.

2.2. Broglia and Winther 1991 (BW 91) potential

BW 91 potential [27] is used as [33]

$$V_{\rm N}^{\rm BW\,91}(r) = -\frac{V_0}{\left[1 + \exp\left(\frac{r - R_0}{a}\right)\right]} \text{ MeV}, \quad (5)$$

where

$$V_0 = 16\pi\frac{R_1 R_2}{R_1 + R_2}\gamma a, \quad a = 0.63\,\text{fm}, \quad (6)$$

and

$$\begin{aligned} R_0 &= R_1 + R_2 + 0.29, \\ R_i &= 1.233A_i^{1/3} - 0.98A_i^{-1/3} \;\; (i = 1, 2), \end{aligned} \quad (7)$$

where γ, γ_0, and k_s, are, respectively,

$$\gamma = \gamma_0\left[1 - k_s\left(\frac{N_p - Z_p}{A_p}\right)\left(\frac{N_t - Z_t}{A_t}\right)\right], \quad (8)$$

$$\gamma_0 = 0.95 \text{ MeV/fm}^2, \;\; k_s = 1.8.$$

2.3. Aage Winther (AW 95) potential

The only difference between AW 95 and BW 91 potentials [33, 34] is

$$a = \left[\frac{1}{1.17(1+0.53(A_1^{-1/3}+A_2^{-1/3}))}\right] \text{fm}, \quad (9)$$

and

$$R_0 = R_1 + R_2, \quad R_i = 1.2A_i^{1/3} - 0.09 \ \ (i = 1, 2). \quad (10)$$

2.4. Bass 1973 (Bass 73) potential

A different version of the proximity potentials from [35, 36] is Bass 73 potential given by [21]

$$V_{\text{N}}^{\text{Bass 73}}(r) = -\frac{da_s A_1^{1/3} A_2^{1/3}}{R_{12}} \exp\left(-\frac{r-R_{12}}{d}\right) \text{MeV}, \quad (11)$$

where

$$R_{12} = 1.07(A_1^{1/3}+A_2^{1/3}), \ \ d = 1.35\,\text{fm}, \ \ a_s = 17\,\text{MeV}. \quad (12)$$

2.5. Bass 1977 (Bass 77) potential

Bass 77 potential [37] is considered as [33]

$$V_{\text{N}}^{\text{Bass 77}}(s) = -\frac{R_1 R_2}{R_1+R_2}\phi(s = r - R_1 - R_2) \ \text{MeV}, \quad (13)$$

where

$$R_i = 1.16A_i^{1/3} - 1.39A_i^{-1/3} \ \ (i = 1, 2). \quad (14)$$

The universal function $\phi(s = r - R_1 - R_2)$ is parametrized by

$$\phi(s) = \left[A\exp\left(\frac{s}{d_1}\right) + B\exp\left(\frac{s}{d_2}\right)\right]^{-1}, \quad (15)$$

where $A = 0.030$ MeV^{-1} fm, $B = 0.0061$ MeV^{-1} fm, $d_1 = 3.30$ fm, and $d_2 = 0.65$ fm.

2.6. Bass 1980 (Bass 80) potential

The difference between Bass 80 and Bass 77 potentials is the function $\phi(s = r - R_1 - R_2)$ shown by [27, 33]

$$\phi(s) = \left[0.033\exp\left(\frac{s}{3.5}\right) + 0.007\exp\left(\frac{s}{0.65}\right)\right]^{-1}, \quad (16)$$

and

$$R_i = R_s\left(1 - \frac{0.98}{R_s^2}\right), \quad R_s = 1.28A_i^{1/3} - 0.76 + 0.8A_i^{-1/3}\,\text{fm} \ \ (i = 1, 2). \quad (17)$$

2.7. Christensen and Winther 1976 (CW 76) potential

CW 76 potential [38] is parametrized by [21]

$$V_{\text{N}}^{\text{CW 76}}(r) = -50\frac{R_1 R_2}{R_1+R_2}\phi(r - R_1 - R_2) \ \text{MeV}, \quad (18)$$

where

$$R_i = 1.233A_i^{1/3} - 0.978A_i^{-1/3} \ \text{fm} \quad (i = 1, 2). \quad (19)$$

The universal function $\phi(s = r - R_1 - R_2)$ is

$$\phi(s) = \exp\left(-\frac{r-R_1-R_2}{0.63}\right). \quad (20)$$

2.8. Ngô 1980 (Ngo 80) potential

The Ngo 80 form of a proximity potential is formulated by [39]

$$V_{\text{N}}^{\text{Ngo 88}}(r) = \overline{R}\phi(r - C_1 - C_2) \ \text{MeV}, \quad (21)$$

$$\overline{R} = \frac{C_1 C_2}{C_1+C_2}, \quad C_i = R_i\left[1 - \left(\frac{b}{R_i}\right)^2 + ...\right], \quad (22)$$

$$R_i = \frac{NR_{ni} + ZR_{pi}}{A_i} \quad (i = 1, 2), \quad (23)$$

$$R_{pi} = r_{0pi}A_i^{1/3}, \ \ R_{ni} = r_{0ni}A_i^{1/3}, \quad (24)$$

$$r_{0pi} = 1.128 \ \text{fm}, \ \ r_{0ni} = 1.1375 + 1.875\times10^{-4}A_i \ \text{fm}. \quad (25)$$

The universal function $\phi(s = r - C_1 - C_2)$ (in MeV/fm) takes the form

$$\Phi(s) = \begin{cases} -33 + 5.4(s-s_0)^2, & \text{for } s < s_0, \\ -33\exp\left[-\frac{1}{5}(s-s_0)^2\right] & \text{for } s \geq s_0, \\ s_0 = -1.6 \ \text{fm}. \end{cases}$$

2.9. Denisov (D) potential

D potential evaluated in the analysis of fusion reactions is given by [33, 40]

$$V_{\text{N}}^{D}(r) = -1.989843\frac{R_1 R_2}{R_1+R_2}\phi(r - R_1 - R_2 - 2.65)\times \times\left[1 + 0.003525139\left(\frac{A_1}{A_2}+\frac{A_2}{A_1}\right)^{3/2} - \right. \left. - 0.4113263(I_1+I_2)\right] \text{MeV}, \quad (26)$$

where

$$I_i = \frac{N_i - Z_i}{A_i}, \tag{27}$$

and

$$R_i = R_{ip}\left(1 - \frac{3.413817}{R_{ip}^2}\right) + 1.284589\left(I_i - \frac{0.4A_i}{A_i + 200}\right), \tag{28}$$

$$R_{ip} = 1.24A_i^{1/3}\left(1 + \frac{1.646}{A_i} - 0.191\left(\frac{A_i - 2Z_i}{A_i}\right)\right) \quad (i = 1, 2). \tag{29}$$

The function $\phi(s = r - R_1 - R_2 - 2.65)$ is considered as

$$\phi(s) = \begin{cases} 1 - \frac{s}{0.7881663} + 1.229218s^2 - 0.2234277s^3 - \\ - 0.1038769s^4 - \frac{R_1R_2}{R_1 + R_2}(0.1844935s^2 + \\ + 0.07570101s^3 + (I_1 + I_2)(0.04470645s^2 + \\ + 0.0334687s^3)) \quad (-5.65 \le s \le 0), \\ \left(1 - s^2\left(0.05410106\frac{R_1R_2}{R_1 + R_2}\exp\left(-\frac{s}{1.76058}\right) - \right.\right. \\ \left.\left. - 0.539542(I_1 + I_2)\exp\left(-\frac{s}{2.424408}\right)\right)\right) \times \\ \times \exp\left(-\frac{s}{0.7881663}\right) \quad (s \ge 0). \end{cases}$$

3. Calculation Procedure

In the present study, the total interaction potential can be assumed as

$$V_{\text{total}}(r) = V_{\text{C}}(r) + V_{\text{N}}(r), \tag{30}$$

where V_{C} is the Coulomb potential shown by [41]

$$V_{\text{C}}(r) = \frac{1}{4\pi\epsilon_\circ}\frac{Z_P Z_T e^2}{r}, \quad r \ge R_c \tag{31}$$

$$= \frac{1}{4\pi\epsilon_\circ}\frac{Z_P Z_T e^2}{2R_c}\left(3 - \frac{r^2}{R_c^2}\right), \quad r < R_c \tag{32}$$

$$R_c = 1.25(A_P^{1/3} + A_T^{1/3}), \tag{33}$$

and V_{N} is the nuclear potential. The real part of the nuclear potential is acquired by using twenty three different versions of a proximity potential. These potentials have been clearly defined in the subsections above. The imaginary part of the nuclear potential is taken as the Woods–Saxon potential shown by

$$W(r) = \frac{W_0}{\left[1 + \exp\left(\frac{r - R_w}{a_w}\right)\right]}, \quad R_w = r_w(A_P^{1/3} + A_T^{1/3}), \tag{34}$$

where W_0 is the potential depth, r_w is the radius parameter, a_w is the diffuseness parameter, and A_P (A_T) is the mass of a projectile (target) nucleus, respectively. The code FRESCO is used in the calculations of fusion cross-sections [42].

4. Results and Discussion

As the first step, we have calculated the nuclear parts of the total interaction potentials of ^{6}He + ^{65}Cuand ^{8}He + ^{65}Cu systems. With this goal, we have achieved the real parts of nuclear potentials, by using twenty three different proximity potentials such as Prox 66, Prox 76, Prox 77, Prox 79, Prox 81-I, Prox 81-II, Prox 81-III, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003-I, Prox 2003-II, Prox 2003-III, Prox 2010, BW 91, AW 95, Bass 73, Bass 77, Bass 80, CW 76, Ngo 80, and D ones. The distance-dependent variations of the real parts of nuclear potentials have been displayed in Fig. 1. Then we have applied the Woods–Saxon potential for the imaginary parts of the nuclear potentials of ^{6}He + ^{65}Cuand ^{8}He + ^{65}Cu reactions. We have set free the W_0, r_w and a_w parameters of the imaginary potential in order to obtain agreement of the theoretical results with the experimental data. We have listed the values of W_0, r_w, and a_w used in the calculations in Tables 2 and 3. When we examine the values of the imaginary potential parameters from Tables 2 and 3, we observe that the imaginary part of the nucleus-nucleus potential in obtaining the cross-sections of fusion reactions has a strong effect on the results.

The fusion cross-sections of ^{6}He + ^{65}Cu reaction have been shown as a function of the center-of-mass energy in Fig. 2. The potential parameters of the imaginary part of the nuclear potential have been given in Table 2. It has been observed that the theoretical results of the proximity potentials except for Bass 73 potential are very similar to each other. The results of other potentials are more smooth and the result of the Bass 73 proximity potential is more oscillating. This makes the result of the Bass 73 potential

even better than other potentials. It can be said that different versions of proximity potentials applied in this work can explain the experimental data on the ^{6}He +^{65}Cu fusion cross-section.

The fusion cross-sections of the ^{8}He +^{65}Cu system have been obtained, by using twenty three different potentials. The results have been compared with the experimental data as a function of the center-of-mass energy in Fig. 3. Additionally, the imaginary potential parameters of the nuclear potential have been listed in Table 3. Similarly to the results of ^{6}He +^{65}Cu reaction, the results of ^{8}He +^{65}Cu reaction except for Bass 73 potential are very close to one another. The results are in a very good agreement with the experimental data. Unlike ^{6}He +^{65}Cu reaction, the results with Bass 73 potential for ^{8}He +^{65}Cu reaction are not better than other potentials. It can be concluded that the proximity potentials evaluated in this study can provide a good agreement with the experimental data on ^{6}He +^{65}Cu fusion cross-section.

Table 2. **The potential parameters W_0 (in MeV), r_w (in fm) and a_w (in fm) used in the calculations of the fusion cross-sections of ^{6}He on ^{65}Cu target nucleus, by using Prox 66, Prox 76, Prox 77, Prox 79, Prox 81-I, Prox 81-II, Prox 81-III, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003-I, Prox 2003-II, Prox 2003-III, Prox 2010, BW 91, AW 95, Bass 73, Bass 77, Bass 80, CW 76, Ngo 80, and D potentials**

Reaction ^{6}He + ^{65}Cu	W (MeV)	r_ω (fm)	a_ω (fm)
Prox 66	15.0	0.9	0.82
Prox 76	15.0	0.9	0.50
Prox 77	15.0	0.9	0.83
Prox 79	15.0	0.9	0.60
Prox 81-I	15.0	0.9	0.70
Prox 81-II	15.0	0.9	0.60
Prox 81-III	15.0	0.9	0.63
Prox 84	15.0	0.9	0.80
Prox 88	25.0	0.9	0.55
Mod-Prox-88	7.00	0.9	0.50
Prox 95	25.0	0.9	0.60
Prox 2003-I	25.0	0.9	0.70
Prox 2003-II	25.0	0.9	0.76
Prox 2003-III	25.0	0.9	0.75
Prox 2010	10.0	0.9	0.50
BW 91	8.80	0.9	0.50
AW 95	5.00	0.9	0.50
Bass 73	4.60	0.9	0.50
Bass 77	19.0	0.9	0.80
Bass 80	4.10	0.9	0.50
CW 76	37.0	0.9	0.50
Ngo 80	17.0	0.9	0.85
D	35.0	0.9	0.87

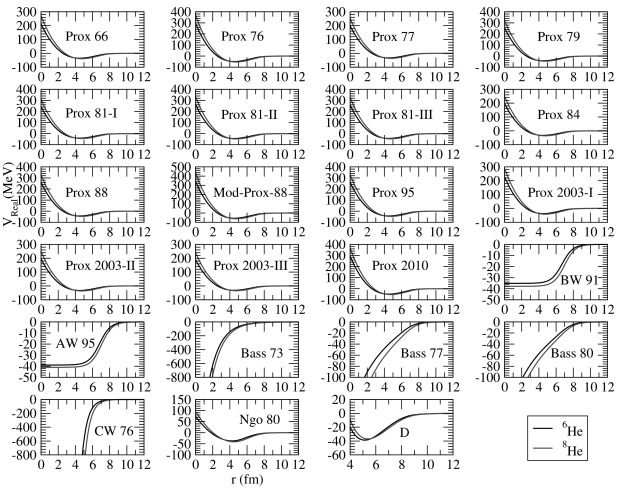

Fig. 1. Distance-dependent changes of the real parts of the nuclear potentials for Prox 66, Prox 76, Prox 77, Prox 79, Prox 81-I, Prox 81-II, Prox 81-III, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003-I, Prox 2003-II, Prox 2003-III, Prox 2010, BW 91, AW 95, Bass 73, Bass 77, Bass 80, CW 76, Ngo 80, and D ones

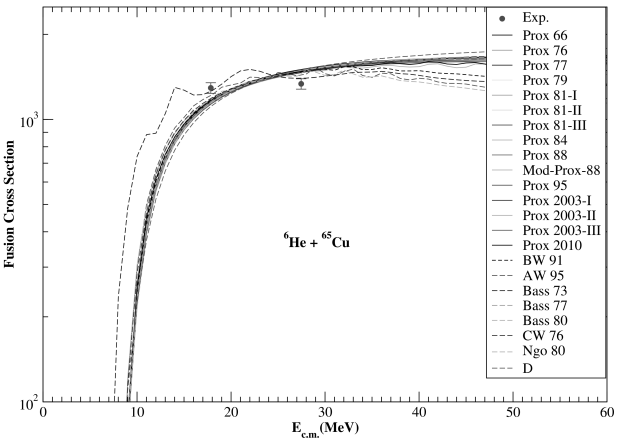

Fig. 2. Fusion cross-sections of ^{6}He + ^{65}Cu reaction in comparison with the experimental data, by using Prox 66, Prox 76, Prox 77, Prox 79, Prox 81-I, Prox 81-II, Prox 81-III, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003-I, Prox 2003-II, Prox 2003-III, Prox 2010, BW 91, AW 95, Bass 73, Bass 77, Bass 80, CW 76, Ngo 80, and D potentials. The experimental data are taken from Ref. [1]

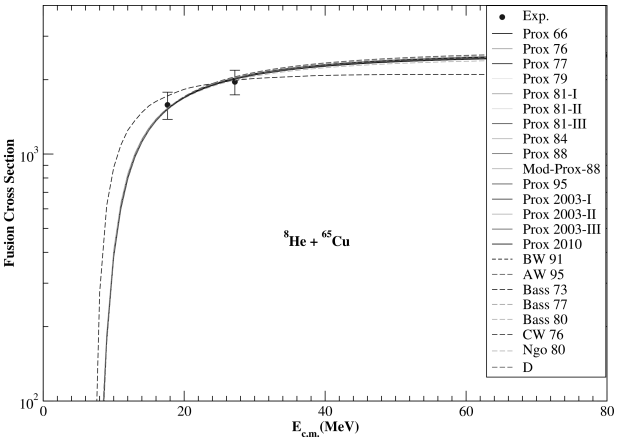

Fig. 3. Fusion cross-sections of ^{8}He + ^{65}Cu reaction in comparison with the experimental data, by using Prox 66, Prox 76, Prox 77, Prox 79, Prox 81-I, Prox 81-II, Prox 81-III, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003-I, Prox 2003-II, Prox 2003-III, Prox 2010, BW 91, AW 95, Bass 73, Bass 77, Bass 80, CW 76, Ngo 80, and D potentials. The experimental data are taken from Ref. [3]

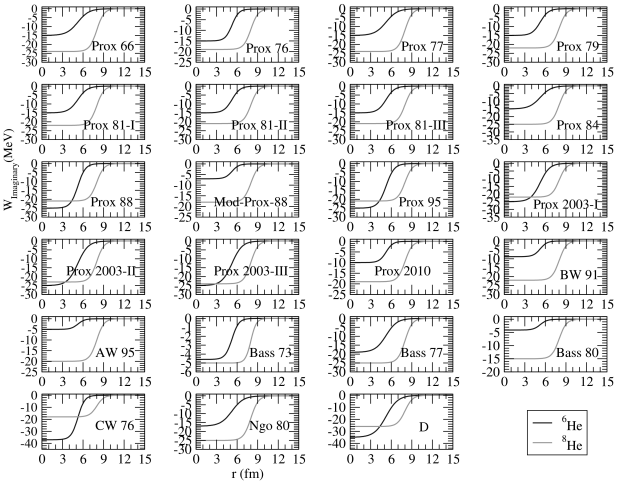

Fig. 4. Distance-dependent changes of the imaginary parts of nuclear potentials for Prox 66, Prox 76, Prox 77, Prox 79, Prox 81-I, Prox 81-II, Prox 81-III, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003-I, Prox 2003-II, Prox 2003-III, Prox 2010, BW 91, AW 95, Bass 73, Bass 77, Bass 80, CW 76, Ngo 80, and D potentials

By carefully examining Figs. 2 and 3, we can notice that the cross-sections of ^{6}He + ^{65}Cu reaction show the appearance of wobbling, while the cross-sections of ^{8}He + ^{65}Cu do not present it. ^{8}He nucleus has a $4n$ + core configuration, and ^{6}He nucleus displays the

Table 3. **The same as in Table 2, but for ^{8}He + ^{65}Cu reaction**

Reaction ^{8}He + ^{65}Cu	W (MeV)	r_ω (fm)	a_ω (fm)
Prox 66	24.0	1.33	0.53
Prox 76	19.0	1.33	0.53
Prox 77	24.0	1.33	0.53
Prox 79	22.0	1.33	0.53
Prox 81-I	22.0	1.33	0.53
Prox 81-II	21.0	1.33	0.53
Prox 81-III	21.0	1.33	0.53
Prox 84	25.0	1.33	0.53
Prox 88	21.0	1.33	0.53
Mod-Prox-88	18.0	1.33	0.53
Prox 95	21.0	1.33	0.53
Prox 2003-I	22.0	1.33	0.53
Prox 2003-II	23.0	1.33	0.53
Prox 2003-III	24.0	1.33	0.53
Prox 2010	19.0	1.33	0.53
BW 91	22.0	1.33	0.53
AW 95	20.0	1.33	0.53
Bass 73	5.00	1.33	0.40
Bass 77	25.0	1.33	0.53
Bass 80	15.0	1.33	0.53
CW 76	18.0	1.33	0.53
Ngo 80	25.0	1.33	0.55
D	26.0	1.33	0.56

$2n$ + core configuration. This leads to the differentiation in the binding energies of ^{6}He and ^{8}He nuclei. As a result, we think that ^{6}He+^{65}Cu reaction displays a wobble structure as compared to ^{8}He + ^{65}Cu reaction. In other words, we can say that the structural differences of ^{6}He and ^{8}He nuclei cause differences in the cross-sections.

The distance-dependent changes of imaginary potentials of ^{6}He and ^{8}He nuclei have been also displayed in Fig. 4. When the depths of the imaginary potentials are examined, it is observed that the potential depth of ^{8}He nucleus is generally deeper than that of ^{6}He nucleus. Additionally, it is seen that ^{6}He nucleus goes to zero faster than ^{8}He nucleus. It can be said that ^{8}He nucleus is more attractive and diffusive than ^{6}He nucleus.

5. Summary and Conclusions

In this study, we have investigated the effects of twenty three different versions of proximity poten-

tials on the fusion cross sections of ^{6}He + ^{65}Cu and ^{8}He + ^{65}Cu reactions. It has been seen that the proximity potentials have given a good agreement results with experimental data. It has been observed that the fusion cross-sections of ^{6}He + ^{65}Cu and ^{8}He + ^{65}Cu systems slightly depend on the shapes of proximity potentials. In addition, it has been noticed that Bass 73 potential in the ^{6}He + ^{65}Cu fusion reaction is slightly better matched to experimental data than other proximity potentials. It has been observed that the results obtained with proximity potentials of ^{8}He + ^{65}Cu fusion reaction, except for Bass 73 potential, are very similar to one another.

Consequently, it can be concluded that different versions of proximity potentials applied in the present work are highly applicable in explaining the experimental data of both ^{6}He + ^{65}Cu and ^{8}He + ^{65}Cu fusion cross-sections. We can say also that the proximity potentials will be interesting in explaining other fusion reactions.

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ВПЛИВ ПОТЕНЦІАЛІВ
БЛИЗЬКОСТІ НА ПЕРЕРІЗ РЕАКЦІЙ
СИНТЕЗУ ^{6,8}He + ^{65}Cu З ГАЛО ЯДРАМИ

Р е з ю м е

Проведено повне теоретичне дослідження потенціалів близькості для найкращого опису реакцій синтезу ^{6,8}He + + ^{65}Cu. Використано такі різні потенціали: Prox 66, Prox 76, Prox 77, Prox 79, Prox 81-I, Prox 81-II, Prox 81-III, Prox 84, Prox 88, Mod-Prox-88, Prox 95, Prox 2003-I, Prox 2003-II, Prox 2003-III, Prox 2010, BW 91, AW 95, Bass 73, Bass 77, Bass 80, CW 76, Ngo 80 і D. Найкращі потенціали визначені на підставі порівняння теоретичних і експериментальних даних по реакціях ^{6,8}He + ^{65}Cu.